Adhesion and friction between individual carbon nanotubes measured using force-versus-distance curves in atomic force microscopy

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(Received 21 April 2008; revised manuscript received 2 June 2008; published 28 July 2008)

The adhesion and friction between individual nanotubes was investigated in ambient using a dynamic atomic force microscope (AFM) operating in force-calibration mode to capture force-versus-distance curves. A multi-walled carbon nanotube (MWNT) tip attached to a conventional AFM probe was brought into contact with and then ramped in vertical direction against a single-walled carbon nanotube (SWNT) bridge suspended over a 2-μm-wide trench. The interaction between nanotubes altered the oscillation amplitude, phase lag, and average deflection of AFM cantilever, from which the interacting forces between nanotubes are quantitatively derived. During ramping, a stick-slip motion was found to dominate the sliding between the nanotubes. The stick was attributed to the presence of high-energy points, such as structural defects or coating of amorphous carbon, on the surface of the MWNT tip. The coefficients of static friction and shear strength between nanotubes were evaluated to be about 0.2 and 1.4 GPa, respectively. They are about 2 orders of magnitude larger than the kinetic counterparts. The kinetic values are on the same order as that measured previously by sliding a MWNT tip across a SWNT bridge in lateral direction.

DOI: 10.1103/PhysRevB.78.045429

PACS number(s): 46.55.+d, 68.35.Np, 61.48.De, 07.79.Lh

I. INTRODUCTION

With their remarkable mechanical and electrical properties,1,2 carbon nanotubes (CNTs) are used to blend with polymers, metals, and ceramics3 or weaved into fibers and sheets4–8 for fabricating macroscale materials with extraordinary strength and toughness and good electrical conductivity. The mechanical strength of these materials critically depends on the mechanical and tribological properties of CNTs, and the small interfacial shear strength between nanotubes and matrix was found to be a limiting factor for enhancing the composite/fiber strength due to inefficiency in interfacial stress transfer.3,7,9,10 To overcome this problem, efforts have been made to introduce defects on the nanotubes by electron irradiation11 or chemical bonds between nanotubes and matrix3,10 in attempts to increase the friction and cross-linking between nanotubes and matrix. However, without quantification of the forces between nanotubes on an individual nanotube level, it is difficult to precisely evaluate how the modification on the structures of individual nanotubes influences the performance of composite/fiber at a macroscale.

The direct nanotribological characterization between two nanotubes is scarce. Very recently our group had addressed this issue by scanning an atomic force microscope (AFM) tip with a multi-walled carbon nanotube (MWNT) welded to its end across a single-walled carbon nanotube (SWNT) bridge suspended on a microtrench in lateral direction.12 The MWNT tip had its outer shells partially stripped possibly with a sword-in-sheath breaking mechanism due to repetitive scans inside the trench,13 leaving a fresh, contamination-free, and defects-minimized surface. The coefficient of kinetic friction between the stripped MWNT tip and the SWNT bridge was estimated to be as low as 0.006, consistent with the value measured on graphite.14,15

Herein, we reinvestigate the friction between individual nanotubes using a pristine MWNT tip freshly grown using catalytic chemical vapor deposition.16 Its surface should bear the original structure as formed in the chemical-vapor deposition (CVD) environment, allowing us to compare with the stripped MWNT tip to gain further insights into the dependence of friction on the surface structure of nanotubes. Also the methodology for measuring the friction between nanotubes is modified to accommodate the increased stiffness of the pristine MWNT tip by operating the AFM in force-calibration mode to capture force-versus-distance curves. In this mode, the MWNT tip is ramped in a vertical direction against a SWNT bridge instead of scanning across it in a lateral direction. The latter is inappropriate for a stiff MWNT tip as it might damage the SWNT bridge during scanning.12 It is well known that on a flat surface, a technique referred to as friction loop measurement is employed as a standard method for measuring friction at the nanoscale using AFM. Here, during the ramping of a MWNT tip in the vertical direction against a SWNT bridge, in the ideal case that the MWNT tip is straight and nondeformable, only the friction would cause the vertical deflection of the AFM cantilever to change. The vertical deflection reverses its sign when the MWNT tip switches its ramping direction from extending to retracting (vice versa) and thus forms a friction loop, too, from which the friction between nanotubes is derived. In view of its similarities with the conventional friction loop measurement, we refer our method as vertical friction loop measurement.

In addition, during the ramping, the AFM cantilever is driven constantly by an embedded piezodriver at its resonance frequency. This enables us to acquire two additional signals, the oscillation amplitude and the phase lag of the AFM cantilever with respect to the driving force. Using a point-mass model to describe the cantilever dynamics, two terms causing the attenuation of amplitude, the interaction stiffness conservative in nature and the damping coefficient dissipative in nature, are extracted from the combination of
amplitude and phase data. By treating the SWNT bridge as an elastic spring with its ends pinned at each edge of the trench, the elastic restoring force and the stiffness of the SWNT bridge are obtained using the strategies developed in Sec. III. The force sensitivity of the AFM probe was then fixed. By manipulating the two parameters, the oscillation amplitude, the horizontal position of the AFM bridge was established, as indicated by an immediate drop of the oscillation amplitude, the horizontal position of the SWNT bridge, and the mechanism leading to this kind of motion is then discussed. Using the simultaneously measured probe is controlled to ramp in the vertical direction (z direction in Fig. 2), while its horizontal position relative to the trench is manipulated by offsetting in the x and y axes. To bring the MWNT tip into contact with the SWNT bridge, the x offset was either gradually increased (the MWNT tip approaches the SWNT bridge from its right side) or decreased (approaching from the left side) in a step of 50 nm, while the oscillation amplitude of AFM cantilever was closely monitored. After contact between the MWNT tip and the SWNT bridge was established, as indicated by an immediate drop of the oscillation amplitude, the horizontal position of the AFM probe was then fixed. By manipulating the two parameters, the oscillation amplitude, the horizontal position of the AFM bridge was established, as indicated by an immediate drop of the oscillation amplitude, the horizontal position of the SWNT bridge, and the mechanism leading to this kind of motion is then discussed. Using the simultaneously measured probe is controlled to ramp in the vertical direction (z direction in Fig. 2), while its horizontal position relative to the trench is manipulated by offsetting in the x and y axes. To bring the MWNT tip into contact with the SWNT bridge, the x offset was either gradually increased (the MWNT tip approaches the SWNT bridge from its right side) or decreased (approaching from the left side) in a step of 50 nm, while the oscillation amplitude of AFM cantilever was closely monitored. After contact between the MWNT tip and the SWNT bridge was established, as indicated by an immediate drop of the oscillation amplitude, the horizontal position of the AFM bridge was established, as indicated by an immediate drop of the oscillation amplitude, the horizontal position of the SWNT bridge, and the mechanism leading to this kind of motion is then discussed. Using the simultaneously measured probe is controlled to ramp in the vertical direction (z direction in Fig. 2), while its horizontal position relative to the trench is manipulated by offsetting in the x and y axes. To bring the MWNT tip into contact with the SWNT bridge, the x offset was either gradually increased (the MWNT tip approaches the SWNT bridge from its right side) or decreased (approaching from the left side) in a step of 50 nm, while the oscillation amplitude of AFM cantilever was closely monitored. After contact between the MWNT tip and the SWNT bridge was established, as indicated by an immediate drop of the oscillation amplitude, the horizontal position of the AFM bridge was established, as indicated by an immediate drop of the oscillation amplitude, the horizontal position of the SWNT bridge, and the mechanism leading to this kind of motion is then discussed. Using the simultaneously measured
oscillation amplitude and phase lag of the AFM cantilever, the conservative and dissipative components in the interaction between nanotubes are separated, and from the latter the average kinetic friction force between nanotubes is derived. A method for determining the contact geometry of nanotubes is introduced. The results are used for evaluating the adhesive and friction force between this MWNT tip and the SWNT bridge.

A. Force-versus-distance curves of MWNT tip ramping against SWNT bridge and the stick-slip motion

The force-versus-distance curves, in which the oscillation amplitude, phase lag and average deflection of AFM cantilever are plotted against its ramping distance, are shown in Fig. 3. They were captured on two consecutive scans due to insufficient data channels available. The MWNT tip was placed below the SWNT bridge and ramped at a height range of about −500 to 300 nm with respect to the trench top. Several feature points are marked in the curve, and for each point, the possible contact geometry between nanotubes is illustrated in the bottom schematics in Fig. 3. During continuous ramping cycles, at point A—the beginning of extending—the SWNT bridge was initially pulled up by the MWNT tip, as the oscillation amplitude of AFM cantilever was significantly lower than its free amplitude, and its deflection was negative (downward). The contact point between nanotubes at point A is denoted by the lower circle drawn onto the MWNT tip in the schematics. With progressive extending of the MWNT tip toward the trench bottom, the restoring force from the SWNT bridge was gradually reduced and then increased when the contact point moved below the trench top. The decrease and increase in the force between nanotubes led to the formation of an arc in the amplitude-distance curve between points A and B. After point B, the restoring force exceeded the static friction between nanotubes, and the SWNT bridge began to slip up along the MWNT tip.

A stick-slip motion of the SWNT bridge was expected from the observed saw-tooth pattern in the amplitude-distance curve between points B and C in Fig. 3. At point C, the MWNT tip was extended to a threshold such that the adhesive force between nanotubes could no longer hold them together, and the two nanotubes detached from each other. The contact point between nanotubes at point of detachment C is denoted by a circle higher than the circle at which the nanotube initially contacted due to the slips between nanotubes. The response of the cantilever here gave a measure of the adhesive force. After detachment, the amplitude was immediately restored to the free amplitude and kept for the rest of the extending from point C to detachment point. At point D, the MWNT tip reversed its ramping direction and moved away from the trench bottom. Contact between the nanotubes appears to be reestablished at point E, as immediate changes to both amplitude and phase are identified. The distance between point C and point E, or the distance between the detachment and attachment of the nanotubes, clearly indicates to what extent the SWNT bridge was displaced. Although the same circle was used to denote the contact point in the schematic, it does not necessarily mean that the contact point at point E is exactly the same as that at point C. The SWNT bridge was pulled with continuous retracting, and at point F, where the contact point should be higher than the trench top, the static friction between nanotubes was overcome, and the SWNT bridge began to slip down along the MWNT tip. Again, a stick-slip motion dominated the sliding between the two nanotubes between points F and A. At the end of the retracting, the contact point slipped to a point denoted by the lower circle in the schematic, where the next ramping cycle started.

Simply by increasing the ramping size by 200 nm while the lowest position of ramping is the same as before, the MWNT tip would break away from the SWNT bridge at the end of retracting and lay on top of it for the following ramping cycles. This allowed us to measure the force-versus-distance curves with the MWNT tip above the SWNT bridge, as shown in Fig. 4. The slight change in contact geometry dramatically altered the way in which the MWNT tip interacted with the SWNT bridge. This is evident by comparing Figs. 3 and 4. In a similar manner, feature points are marked on the curve, and corresponding contact geometries are illustrated. One can refer to the schematic in Fig. 4 for an intuitive picture of the ramping process. A few points worthy of mention are as follows. (1) A compressive force between the nanotubes was developed between points C and D in Fig. 4.
coating of amorphous carbon. The shape of the MWNT tip grown in CVD environment are rich in structural defects and elastic deformation at those points. The arrows in the schematics indicate the ramping direction at those points. The contact geometries of nanotubes are illustrated in the schematics from above. Feature points are labeled in the top plots, and the MWNT tip ramping in vertical direction against a SWNT bridge is shown in Fig. 1. During the ramping of the MWNT tip, the SWNT bridge tended to stick to a HEP unilaterally, instead of a pulling force in the region between point B and C in Fig. 3. Thus, the force worked collaboratively with the adhesive force between nanotubes to enhance the friction between them instead of weakening it in Fig. 3. (2) The detachment of nanotubes happened at the end of retracting in Fig. 4 instead of the end of extending in Fig. 3, where the contact points between nanotubes are expected to be different. (3) The saw-tooth pattern observed in Fig. 4 is significantly different from that in Fig. 3. In addition to different load between the nanotubes, the different sides which the MWNT tip used to contact the SWNT bridge are likely to be another contributor, indicating the stick-slip motion is surface sensitive.

The lack of atomic scale periodicity in the stick-slip motion (Figs. 3 and 4) suggests its origin from random high-energy points (HEPs) presented on the surface of nanotubes, most likely on the MWNT tip, rather than the periodic atomic lattice found previously. It is known that MWNTs grown in CVD environment are rich in structural defects and coating of amorphous carbon. The shape of the MWNT tip used here is far from the perfect cylinder in the SEM image of the MWNT tip shown in Fig. 1. During the ramping of the MWNT tip, the SWNT bridge tended to stick to a HEP unless the restoring force of the SWNT bridge grew strong enough to shear it away. After that the SWNT bridge slipped along the MWNT tip under the effect of the restoring force, the MWNT tip used to contact the SWNT bridge are likely to be different, preventing an accurate estimation of the lower limit of the static friction force, considering the slippage between nanotubes could happen at point F in Fig. 3 where the amplitude is close to the free amplitude.

In our case, the average distance between HEPs appeared to be smaller than the displacement of SWNT bridge, which led to resticking of the SWNT bridge to subsequent HEPs before the strain of SWNT bridge decreased to zero. The sticking force, i.e., static friction, between nanotubes was large enough to be detected as a variation of cantilever deflection (Figs. 3 and 4). In the ideal case that the MWNT tip is straight and nondeformable, only the friction force between nanotubes will contribute to the vertical deflection, as shown in Fig. 2. However, in reality, the MWNT tip here was tilted toward the −x direction and could be deformed nonlinearly, and thus the vertical deflection always involved contributions from both normal and friction forces.

Figure 5 shows possible contact geometry with the MWNT tip below the SWNT bridge. The angles between the MWNT tip and the SWNT bridge are α and β, respectively. The friction force $F_F$ is related to the vertical force $F_V$, which is an experimentally measurable quantity, by a conversion factor, $\cos \beta/\cos \alpha$. Generally speaking, to precisely calculate the friction force, one needs to determine these two angles. Later we will introduce a method for the evaluation of one of the angles, $\alpha$. At present we simply utilize the fact that when the contact point between nanotubes is much higher than the trench top, these two angles should be close to zero, leading to a conversion factor of about 1. This means that we can use the cantilever deflection data from the right portion of force-versus-distance curves, where the contact point is higher than the trench top to calculate the friction force between nanotubes. Seventeen force-versus-distance curves captured with the MWNT tip below the SWNT bridge are summarized in Fig. 6(a) by plotting the vertical force (normal spring constant × vertical deflection) versus the normalized amplitude. The left and right portions of force-versus-distance curves split into the upper and lower branches in Fig. 6(a) with the left portion corresponding to the upper branch and vice versa. The static friction force estimated from the lower branch varied between 0 and 1.3 nN. The large noise in the deflection data, including thermal noise on the order of 0.1 nN (Ref. 12) and strong laser interference, prevents an accurate estimation of the lower limit of the static friction force, considering the slippage between nanotubes could happen at point F in Fig. 3 where the amplitude is close to the free amplitude.
FIG. 6. (Color online) The interaction between nanotubes is partially detected as vertical deflection of AFM cantilever, from which the vertical force is calculated and plotted against normalized amplitude using (a) 17 force-versus-distance curves with the MWNT tip below the SWNT bridge and (b) 12 force-versus-distance curves with the MWNT tip above the SWNT bridge.

The detachment of the nanotubes near the end of extending in Fig. 3 and the end of retracting in Fig. 4 allows estimation of the adhesive force between nanotubes at different contact points. In Fig. 6(a), 12 force-versus-distance curves captured with the MWNT tip above the SWNT bridge are also summarized. The vertical forces are almost the same and close to 3.0 nN for the highest point in Fig. 6(a) and the lowest point in Fig. 6(b). This is an unexpected result since the conversion factor for the adhesive force from the vertical force, $k_z$, is the spring constant of the cantilever, and $A$ and $\varphi$ are the attenuated oscillation amplitude and phase lag of AFM cantilever. The normalized damping coefficient is given by

$$
\frac{k_{\text{int}}}{k_z} = \left( \frac{\omega}{\omega_0} \right)^2 - 1 + \frac{\cos \varphi + A_\varphi/(A \sqrt{1 + Q_0^2})}{A \sqrt{1 + Q_0^2}/A_0 + 2 \cos \varphi + A_\varphi/(A \sqrt{1 + Q_0^2})},
$$

where $k_z$ is the spring constant of the cantilever, and $A$ and $\varphi$ are the attenuated oscillation amplitude and phase lag of AFM cantilever. The normalized damping coefficient is given by

$$
\frac{\gamma_{\text{int}}}{\gamma_0} = -\frac{\omega_0}{\omega} A \sqrt{1 + Q_0^2}/A_0 + 2 \cos \varphi + A_\varphi/(A \sqrt{1 + Q_0^2}) - 1,
$$

where $\gamma_0$ is the bulking damping coefficient in air.

Using Eqs. (1) and (2), the normalized interaction stiffness and the damping coefficient are calculated from various combinations of normalized oscillation amplitude $A/A_0$ and phase lag $\varphi$ and plotted as contour plots in Fig. 7. The measured amplitude and phase with the MWNT tip below and above the SWNT bridge are mapped into the contour plots for straightforward evaluations of the interaction stiffness and damping coefficient. It can be concluded from Fig. 7 that the interaction between nanotubes is dominated by the conservative term, the interaction stiffness. During the interaction of nanotubes, the SWNT bridge acted as an elastic spring, and its stiffness was added to the AFM cantilever through the MWNT tip. The dissipative term is small as all data basically follow the zero contour line of damping coefficient. The deviation of the bottom data from the zero contour line should be an artifact associated with the inaccurate measurements of amplitude and phase data at small amplitude. Moreover, the difference between the extending and retracting data is small, although load and contact geometries are expected to be significantly different, suggesting the deviation results from measurement errors. The domination of interaction stiffness in the interaction between nanotubes confirms the stick-slip motion as the sliding mechanism be-
between this MWNT tip and the SWNT bridge since only during sticking could the stiffness of SWNT bridge be effectively transferred to the AFM cantilever.

The average kinetic friction between nanotubes is estimated from the amplitude and phase data using a formula developed by Bhushan et al.,\textsuperscript{12}

\[ F_F = \frac{\pi k_v (A_0 \sin \varphi - A)}{4Q} \]  

It can also be obtained by integrating the instantaneous power dissipation over one oscillation period from Eq. (2) with an approximation that \( \gamma_{in} / \gamma_0 = -A_0 \sin \varphi / A - 1 \) under high \( Q_0 \). The results are plotted against normalized amplitude in Fig. 8. Except for a few transient points, the kinetic friction force never exceeded 20 pN. The kinetic friction force is on the same order as the value previously measured by Bhushan et al.\textsuperscript{12} by sliding a MWNT tip across a SWNT bridge in lateral direction, as expected for shearing an incommensurate contact between graphite sheets.\textsuperscript{23} Note that unphysical negative frictions are generated at small amplitude due to inaccurately measured amplitude and phase data. However, the difference between the forces during the extending and retracting is still within 20 pN for the data at small amplitude.

\[ k_{sw} = \frac{4 \pi k_v}{\gamma_0} \left[ 1 - \frac{1}{(1 + \tilde{d})^{3/2}} \right] \]  

The value of \( F_v / k_{sw}^2 \) is a function of \( \tilde{d} \) only. For each data point in Figs. 3 and 4, \( F_v \) is estimated directly from the measured vertical deflection, and \( k_{sw} \) is estimated from the amplitude by assuming \( \gamma_{in} = 0 \) using Eqs. (1) and (2). The displacement is then calculated numerically from the value

FIG. 7. (Color online) The relative magnitude of interaction stiffness and damping coefficient for the interaction between nanotubes are determined by mapping (a) the data in the top figures of Fig. 3 and (b) in the top figures of Fig. 4 into the contour plots of interaction stiffness and damping coefficient against normalized amplitude and phase. The solid lines represent the damping coefficient and the dashed lines represent the interaction stiffness. The data basically follow the zero contour line of damping coefficient highlighted by a thick solid line.

C. Determination of the contact geometry for calculating the adhesive force between nanotubes

This section is devoted to the determination of the displacement of the SWNT bridge (Fig. 5), from which the conversion factor to obtain the adhesive force from the vertical force can be calculated. Due to the complexity involved in the three-dimensional contact between nanotubes, a few simplifications need to be made before the displacement of the SWNT bridge becomes computable with present data.

First, the SWNT bridge is assumed to be perpendicular to the trench wall, and the MWNT tip is not tilted in the y direction. Second, the MWNT tip is assumed to contact the SWNT bridge at the center of the bridge. These two assumptions are necessary to reduce the contact problem to two dimensions, i.e., the \( xz \) plane. Third, the MWNT tip is assumed to be much stiffer than the SWNT bridge so that its contribution to the interaction stiffness can be neglected. With these assumptions, the unknown quantities, restoring force and stiffness of the SWNT bridge, are directly related to the measurable quantities, vertical force and interaction stiffness, through the displacement of the SWNT bridge. Referring to Fig. 5, the interaction stiffness is a component of the stiffness of the SWNT bridge, \( k_{sw} = k \cos^2 \alpha \), and the vertical force is a component of the restoring force, \( F_v = F \cos \alpha \). By treating the SWNT bridge as an elastic string, the restoring force is a simple function of the displacement of the SWNT bridge, derived from the equations given by Walters et al.,\textsuperscript{23}

\[ F = 2EA \left( \frac{\bar{d} - \tilde{d}}{\sqrt{\bar{d}^2 + 1}} \right) \]  

where \( A \) and \( \bar{d} \) are the cross-sectional area and the normalized displacement of SWNT bridge (\( \bar{d} = 2d/g \), where \( g \) is the gap width of the trench). Its derivative gives the stiffness for the SWNT bridge,

\[ k = \frac{4EA}{g} \left[ 1 - \frac{1}{(1 + \tilde{d})^{3/2}} \right] \]  

It can also be obtained by integrating the instantaneous power dissipation over one oscillation period from Eq. (2) with an approximation that \( \gamma_{in} / \gamma_0 = 0 \) using Eqs. (1) and (2).
FIG. 9. (Color online) The contact geometry between nanotubes is partially determined and the distance (left column) and angle (right column) of the displacement of SWNT bridge are plotted against normalized amplitude. The plots are calculated from (a) Fig. 3 and (b) Fig. 4. The results are used for precise evaluation of the adhesive and friction force between nanotubes.

D. Determination of the adhesive force, friction coefficient, and shear strength between nanotubes

The adhesive force between nanotubes is related to the vertical force by a conversion factor of $\sin \beta / \cos \alpha$. With the MWNT tip below the SWNT bridge, the detachment of nanotubes occurred near the end of extending at point C in Fig. 3. At point C, the MWNT tip was pulled up by the SWNT bridge, and the angle between the nanotubes, $\beta$, would increase with this geometry. Also the adhesive force between nanotubes is counterbalanced by the pulling force from the SWNT bridge (which is the reason the nanotubes detached from each other), leading to an almost zero friction force at point C. However, the vertical force is large at this point, close to 3.0 nN, which means that the angle $\beta$ must be close to 90° to minimize its shear component, which is balanced by the nearly zero friction force. In addition, we estimate $\cos \alpha$ of 0.5 for data at point C, which leads to a conversion factor of 2, and thus the adhesive force between nanotubes should be 6.0 nN from the measured vertical force of 3.0 nN. In Fig. 9(b), the MWNT tip detached the SWNT bridge from above, and the angles $\alpha$ and $\beta$ would be both close to zero. Thus the adhesive force between nanotubes should be smaller than the measured vertical force of 3.0 nN for a conversion factor smaller than 1. The difference between the adhesive forces indicates that the MWNT tip might detach from the SWNT bridge using its tip end when it was placed above the SWNT bridge instead of using its shell when it was placed below the SWNT bridge. The adhesive force between the end of the MWNT tip and the shell of SWNT bridge was reported to be 0.7 nN by Bhushan et al.,12 about an order of magnitude smaller than the adhesive force between the shells of the nanotubes measured here, possibly due to the small radius of curvature at the MWNT tip end.

By comparing the friction force to the adhesive force, an upper limit for the coefficient of static friction is estimated to be 0.2, about 2 orders of magnitude larger than the coefficient of kinetic friction of 0.003. The latter is on the same order as the value reported previously for the stripped MWNT tip.12 The HEPs, which led to the 1 nN static friction during the sliding between nanotubes, were reported to increase the interlayer sliding friction inside MWNT to the 1 nN range as well.24 While the HEPs in the interlayer sliding experiments could be dynamic due to e-beam radiation,24 the...
HEPs on this MWNT tip appear to be intrinsic and permanent as the repeatability of force-versus-distance curves were reasonably good. This finding highlights the importance of the surface structure of nanotubes on their nanotribological properties.

The interface energy and shear strength between nanotubes are estimated using the Johnson-Kendall-Roberts model. The MWNT has a diameter of 70 nm, and the SWNT has an average diameter of 1.43 nm. The measured adhesion force of 6.0 nN corresponds to a work of adhesion between nanotubes of 0.25 N/m, as estimated using the equation

\[ W_A = \frac{2F_A}{3\pi\sqrt{r_1r_2}}, \]

where \( r_1 \) and \( r_2 \) are the radii of the MWNT tip and the SWNT bridge, respectively. This value agrees with the cohesive energy of graphite measured using other methods, justifying the use of the contact geometry estimation method. The contact area between nanotubes at zero loads is given by

\[ A = \pi(6\pi W_A r_1 r_2/K)^{2/3}, \]

where \( K \) is the composite Young’s modulus, given as \( K = 2E/(3(1-v^2)) \), where \( v \) is the Poisson ratio of nanotube, assumed to be 0.3. Using a reported value of 0.22 N/m as the work of adhesion between nanotubes, a contact area at zero force of 0.85 nm\(^2\) is obtained. The shear strength obtained by dividing the static friction with the contact area could be as high as 1.4 GPa, which is about 2 orders of magnitude larger than the shear strength calculated using kinetic friction.

**IV. CONCLUSION**

In this paper, a method, referred to here as vertical friction loop measurement, was proposed to characterize the adhesion and friction properties between individual carbon nanotubes using the force-calibration-mode AFM in ambient. A MWNT tip attached to a conventional AFM probe was ramped in the vertical direction against a SWNT suspended on a microtrench at crossed geometry. The adhesion and friction between nanotubes were evaluated from the response of AFM cantilever during ramping.

During ramping, a stick-slip motion was found to dominate the sliding between the nanotubes. We attribute the stick to the presence of HEPs on the surface of this MWNT tip, associated with structural defects or coating of amorphous carbon. The stick was broken once the restoring force of the SWNT bridge became large enough to overcome the static friction between nanotubes, and the nanotubes began to slip against each other. The slip events were transient due to the fast shifting of the SWNT bridge between the HEPs separated by small distance.

The dynamics of the AFM cantilever were analyzed using a point-mass model, and the amplitude attenuation of the AFM cantilever during the interaction of nanotubes was found to be primarily attributed to the increase in the effective cantilever stiffness. The dissipative contribution from the kinetic friction between nanotubes to the amplitude attenuation is significantly smaller than the contribution from the conservative interaction stiffness.

The SWNT bridge acts as an elastic spring; and its restoring force and stiffness are functions of its displacement and related to the measured vertical force and interaction stiffness. By analyzing their relations, the distance and angle of displacement were calculated. They are used to calculate the adhesive and friction forces between nanotubes by taking into account the contact geometry between nanotubes.

The intershell adhesion force was found to be 6.0 nN, from which the work of adhesion between nanotubes is estimated to be 0.25 N/m, consistent with the value measured using different methods. The static friction was estimated to be as high as 1.3 nN from the vertical deflection of cantilever during the stick-slip motion. Dividing the static friction by the adhesive force yields a coefficient of static friction of 0.2, and the shear strength between nanotubes evaluated from the static friction is 1.4 GPa. The average kinetic friction between nanotubes was estimated to be smaller than 20 pN, which leads to a kinetic friction coefficient of 0.003, comparable to the value reported before. The coefficient of static friction is about 2 orders of magnitude larger than that of the coefficient of kinetic friction, emphasizing the importance of the surface structure of nanotubes on their nanotribological properties.

**ACKNOWLEDGMENTS**

The authors gratefully acknowledge Alain Jungen for providing the SWNT sample, Cattien V. Nguyen for providing the MWNT probe, and Kwang Joo Kwak for SEM measurements.

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PHYSICAL REVIEW B 78, 045429 (2008)